

1. Physics 13: Electrical and Thermal Properties -Metals

Physics 13: Electrical and Thermal Properties - Metals

- Classical free electron theory - Ohm's Law
  - But fails: magnitude of  $\rho$  (resistivity) or  $\sigma$  (conductivity)
  - $\rho_{\text{classical}}(T) \sim v_{\text{ave}}(T) \sim \sqrt{T}$  but  $\rho_{\text{msd}}(T) \sim T$  &  $\rho_{\text{msd}} \sim \rho_{\text{class}}/10$ 
    - QM electron gas:  $v_{\text{ave}} \sim$  independent of T (for  $kT \ll E_{\text{fermi}}$ )
    - $\rho$  depends on wave-like  $e^-$ s propagating in lattice
  - Why conductors and insulators?
  - Heat conduction and capacity are incorrect
  - Classical  $e^-$  gas:  $eE - F_{\text{drag}} = 0$  &  $F_{\text{drag}} \sim v_{\text{drift}} / (\text{collision time } \tau)$  among ions. So  $v_{\text{drift}} \sim E \tau$ . Current  $j \sim nev_{\text{drift}}$ . Ohm's law & resistivity  $\rho \sim v_{\text{drift}} \sim 1/\tau$ . But  $\tau \sim 1/v_{\text{AVE}}$  so  $\rho \sim v_{\text{AVE}}$  & from stat.mech.  $v_{\text{AVE}} \sim \sqrt{T}$ . What is wrong with this picture?

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2. Electrical & Thermal properties - QM

Electrical & Thermal properties - QM

- $e^-$ s obey exclusion principle - occupy states of solid lattice accordingly - even at  $T=0^{\circ}\text{K}$  they have finite  $E$ 's
- 1-dim Square well model at  $0^{\circ}\text{K}$ 
  - $E_1 = h^2/8m_e L^2$  & for N  $e^-$ s each state ( $E_n = n^2 E_1$ ) is occupied by a pair
  - Highest E for  $N/2$  or  $E_{\text{Fermi}}(T=0K) = E_{N/2} = N^2 E_1/4$

$$\text{or } E_{\text{Fermi}}(T=0K) = \frac{h^2 c^2}{32 m_e c^2} \left(\frac{N}{L}\right)^2 = \frac{(1240 \text{ eV} \cdot \text{nm})^2}{32 (5.11 \times 10^5 \text{ eV})} \left(\frac{N}{L}\right)^2 = 9.40 \times 10^{-2} \text{ eV} \cdot \text{nm}^2 \left(\frac{N}{L}\right)^2$$

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3. 1-dim Fermi level example - Cu

### 1-dim Fermi level example - Cu

Cu has  $\sim 1$   $e^-$  per atom. Density =  $8.92 \text{ g/cm}^3$  (tables).  
 Molar mass =  $63.5 \text{ g/mole}$ . So  $8.92/63.5 = 0.140 \text{ mole/cm}^3$ .  
 Then  $N/V = 0.140 \times (6.02 \times 10^{23}) = 8.43 \times 10^{22} \text{ e/cm}^3$ .  
 Taking a cubic lattice to get 1-d  
 $N/L = (8.43 \times 10^{22})^{1/3} = 4.39 \times 10^7 / \text{cm} = 4.39 / \text{nm}$ .  
 So 1-d Fermi level is  
 $E_{\text{Fermi}} = 9.4 \times 10^{-2} \times (N/L)^2 = 9.4 \times 10^{-2} \times (4.39)^2 = 1.8 \text{ eV}$  (?)  
 This 1-dim model result  $\rightarrow$  right order of magnitude  
 Compare that to  $kT(300\text{K}) = 2.6 \times 10^{-2} \text{ eV}$ .

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4. 3-dim Fermi level

### 3-dim Fermi level

- In 3-d Fermi energy is altered -  $e^-$ s now fill a spherical well (real potentials?)

$$E_{\text{Fermi}}(T = 0\text{K}) = \frac{h^2 c^2}{8m_e c^2} \left( \frac{3N}{\pi V} \right)^{2/3}$$

Examples:	$N/V(\text{cm}^3)$	$E_F(\text{eV})$
Al	$18.1 \times 10^{22}$	11.7
Cu	$8.47 \times 10^{22}$	7.04
K	$1.40 \times 10^{22}$	2.11

$E_F(0\text{K}) \gg kT(300\text{K})$   
(or 0.026eV)

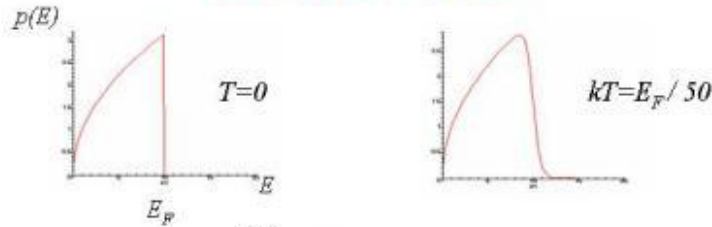
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5. Density of states

Density of states



At  $T=0K$   $p(E) = \frac{3N}{2} E_F^{-3/2} E^{1/2}$  for  $E \leq E_F$  and 0 otherwise.

$$\int_0^{E_F} p(E) dE = N \quad \text{and} \quad E_{ave} = \int_0^{E_F} E p(E) dE / N = \frac{3}{5} E_F$$

At finite T have Fermi - Dirac factor  $\frac{1}{e^{\frac{E-E_F}{kT}} + 1}$  also

Only small fraction of e's ( $\sim kT/2E_F$ ) above  $E_F$  can move to other states in response to Electric field.

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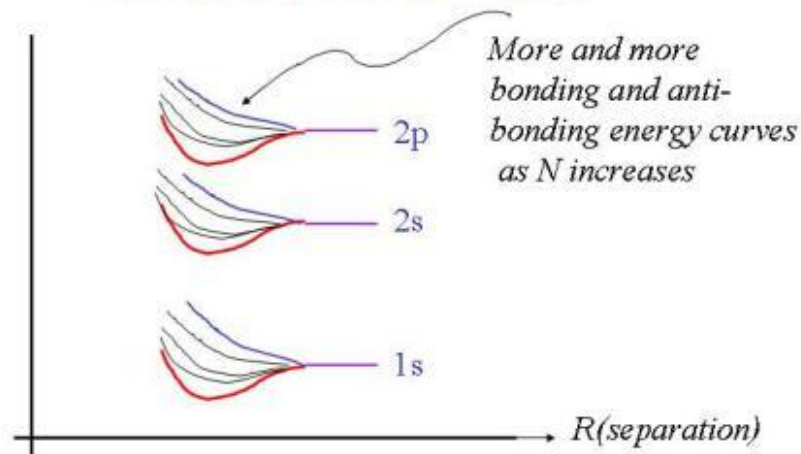
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6. Filling energy states for lattice of atoms (beyond square we...

Filling energy states for lattice of atoms  
(beyond square well model)



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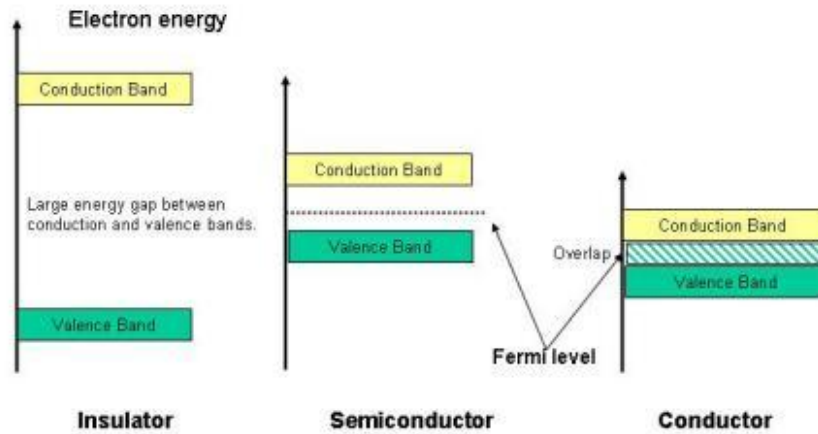
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7. Band theory of conductivity

### Band theory of conductivity



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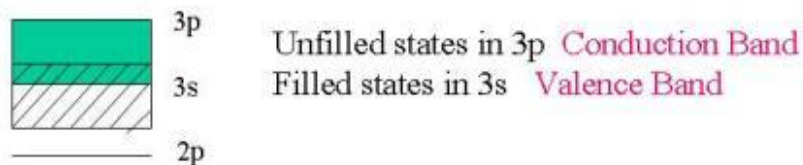
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8. Examples - Na, Mg - Conductors

### Examples - Na, Mg - Conductors

- Na atom K & L shells filled +  $3s^1$  valence  $e$ 's
- In solid 1/2 of **3s Band** is filled
  - 3s electrons are free to move, i.e. jump to empty states  $\Rightarrow$  conductor
- Mg: atomic  $3s^2$  valence  $e$ 's
- In solid 3s Band has some overlap with 3p Band  $\Rightarrow$  conductor



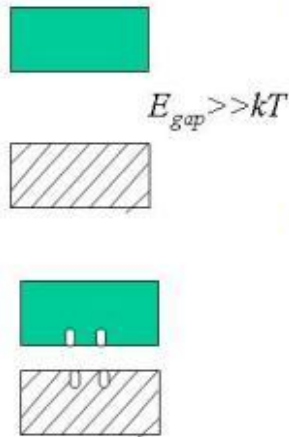
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9. Insulators & semiconductors

Insulators & semiconductors



- Conduction band empty  
Valence band full  
Big gap  $\Rightarrow$  **Insulator**  
(e.g. NaCl, C)
- Small gap  $\sim kT$  - some e's excited to conduction band at T - can be accelerated by E - holes also - **Semiconductor** (Si)
- Intrinsic & doped(n and p)

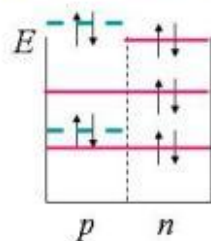
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10. Nuclear properties

Nuclear properties

- Z=atomic number: Nuclide  
A=N+Z= mass no.: Isotope ( ${}^1_1\text{H}$ ,  ${}^2_1\text{H}=\text{d}$ ,  ${}^3_1\text{H}=\text{t}$ )  
 $m_n > m_p$  (0.2%) : Nucleons
- Size:  $R=R_0 A^{1/3}$  ( $R_0 \sim 1.5 \text{ fm}$ ,  $\text{fm}=\text{femtometer}=10^{-15}\text{m}$ )
- Strong (or Hadronic) Force  $\sim$  same for n & p
- $N \geq Z$  : Coulomb repulsion elevates proton levels



n&p are bound in Strong Nuclear Potential  
Nuclear quantum states and  $\Psi$   
For higher Z next n levels will be considerably lower in E than next p levels - e.g.  ${}_{92}^{238}\text{U}$

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## Nuclear properties II

### Nuclear properties II

- Binding Energy  $E_{BE} = (Zm_p + Nm_n - M_A)c^2$ 
  - Mass defect, curve of binding energy, fission, fusion
- Magnetic moments: nucleons first

$$\mu_{Nuclear} = \frac{e\hbar}{2m_p} = 3.15 \times 10^{-8} \text{ eV / Tesla (cf. } 5.79 \times 10^{-5} \text{)}$$

$$\mu_z^p = +2.79\mu_N \quad \mu_z^n = -1.91\mu_N$$

- Nuclear Magnetic Moments - combinations
- Nuclear Magnetic Resonance and MRI scans

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## Nuclear decays

### Nuclear decays

- $\alpha$ :  $(Z,A) \rightarrow (Z-2,A-4) + {}_2^4\text{He}$ 
  - Strong or nuclear glue and quantum tunneling
- $\beta$ :  $(Z,A) \rightarrow (Z+1,A) + e^-$  also have  $(Z-1,A) + e^+$ 
  - Weak and based on  $n \rightarrow p + e^- + \bar{\nu}$  (*anti-neutrino*)
- $\gamma$ :  $(Z,A)^* \rightarrow (Z,A) + \gamma$ 
  - EM quantum transition
- Each with characteristic decay time and energy release (Q value)

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13. Radioactivity

### Radioactivity

- All decays occur randomly - probability distribution characteristic of nuclide and decay products
- Activity=rate of decrease in time= $-dN/dt=A(t)$
- $A$  in units of Curies= $3.7 \times 10^{10}$  disintegrations/sec or Becquerels (Bq) SI units
- $A \propto N(t)$  (each decay is independent of others) with  $A = \lambda N(t)$   $\lambda$ =decay constant
- How does  $N$  change in time?

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14. Exponential decay

### Exponential decay

$$A(t) = -\frac{dN}{dt} = \lambda N(t) \quad \text{So} \quad \frac{dN}{N} = -\lambda dt$$

$$\text{and} \quad \int_{N_0}^{N(t)} \frac{dN}{N} = -\int_0^t \lambda dt \quad \text{or} \quad \ln(N(t)) - \ln(N_0) = -\lambda t$$

$$\text{or} \quad N(t) = N_0 e^{-\lambda t} \quad \text{and} \quad A(t) = A_0 e^{-\lambda t}$$

$$\text{Hence} \quad \frac{1}{\lambda} = \tau \quad \text{mean lifetime}$$

$$\text{Note also that } A(t) \text{ will be } \frac{1}{2}A_0 \text{ for } e^{-\lambda t} = \frac{1}{2} \text{ or } t_{1/2} = \frac{\ln 2}{\lambda}$$

$t_{1/2}$  is the half-life. Exponential decay law for many phenomena, especially QM decays and heat transfer.

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## Decay examples

### Decay examples

• NUCLEUS	ATOMIC WEIGHT	HALF LIFE	DECAY PRODUCTS
• Ra	226	1602 years	alpha,gamma(4%)
• C	14	5730 years	beta
• Sr	90	28 years	beta
• <sup>3</sup> H	3	12.3 yr	beta
• I	131	8 days	beta,gamma
• Cs	137	30 yr	beta,gamma
• Rn	222	3.8 days	alpha,gamma(0.1%)
• U	238	4.50 x 10 <sup>9</sup> yr	alpha,gamma(23%)
• Pu	239	24,400 yr	alpha

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